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Electronic sideband locking of a broadly tunable 318.6nm ultraviolet laser to an ultra-stable optical cavity

Jiandong Bai¹,², Jieying Wang¹,², Jun He¹,²,³ and Junmin Wang¹,²,³

¹ State Key Laboratory of Quantum Optics and Quantum Optics Devices, Shanxi University, Tai Yuan 030006, Shan Xi Province, People’s Republic of China
² Institute of Opto-Electronics, Shanxi University, Tai Yuan 030006, Shan Xi Province, People’s Republic of China
³ Collaborative Innovation Center of Extreme Optics, Shanxi University, Tai Yuan 030006, Shan Xi Province, People’s Republic of China

E-mail: wwjjmm@sxu.edu.cn

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Abstract
We demonstrate frequency stabilization of a tunable 318.6 nm ultraviolet (UV) laser system using electronic sideband locking. By indirectly changing the frequency of a broadband electro-optic phase modulator, the laser can be continuously tuned over 4 GHz, while a 637.2 nm laser is directly stabilized to a high-finesse ultra-stable optical cavity. The doubling cavity also remains locked to the 637.2 nm light. We show that the tuning range depends mainly on the gain-flattening region of the modulator and the piezo-tunable range of the seed laser. The frequency-stabilized tunable UV laser system is able to compensate for the offset between reference and target frequencies, and has potential applications in precision spectroscopy of cold atoms.

Keywords: electronic sideband locking, frequency offset locking, ultraviolet laser, continuously tunable laser frequency

(Some figures may appear in colour only in the online journal)

1. Introduction

Strong long-range dipole–dipole interactions between highly excited Rydberg atoms lead to the Rydberg blockade mechanism [1–3], which has been used in many practical quantum-based technologies, such as logic gates [4], entanglement between two atoms [5, 6], communications [7] and information processing [1]. In most situations, a two- or three-step excitation is used to address a specific Rydberg state. However, this multi-step excitation inevitably populates intermediate states, producing photon scattering and ac-Stark shifts resulting in low excitation efficiency. These disadvantages can be avoided with single-photon excitation. Tong et al [8] and Hankin et al [9] have recently demonstrated direct single-photon Rydberg excitation of ⁸⁷Rb at 297 nm and ¹³³Cs at 319 nm. For single-photon Rydberg excitation of cesium atoms from the 6S₁/₂ to an nP (n = 70–100) state, a 318.6 nm ultraviolet (UV) laser should have high power, frequency stabilization, and be continuously tunable over a wide range. Previously [10], we followed the pioneering work of Wineland et al [11] to generate a 2.26 W 318.6 nm UV laser starting from two infrared fiber lasers.

A common method to produce a frequency-stabilized tunable laser is to lock it to a fixed optical cavity, or to specific atomic or molecular spectral lines, and then insert one or more acousto-optic modulators (AOMs) in the optical path [12]. A given AOM can create MHz-level shifts, but has a limited bandwidth of dozens of MHz that is only a small fraction of its fixed center frequency. Special AOMs that generate shifts of GHz offer an increased bandwidth (less than several hundred MHz), but their diffraction efficiency is normally very low, and they are very expensive. For a bulk-
type electro-optic phase modulator (EOPM), its maximum bandwidth is only hundreds of MHz. In addition, they need large radio-frequency power consumption. Compared with the bulk-type EOPM and AOM, commercially available fiber-based waveguide-type EOPMs avoid these limitations and have typical bandwidths up to dozens of GHz. Thorpe et al [13] proposed sideband locking techniques that involve a few modifications of standard Pound–Drever–Hall (PDH) [14] locking. Based on a waveguide-type EOPM, serrodyne sideband modulation [15, 16] has been recently used for the frequency stabilization of a tunable laser to an optical cavity, with a typical dynamic range of 220 MHz.

Here, we use the electronic sideband (ESB) locking scheme and a high-finesse ultrafast expansion (ULE) optical cavity to produce a frequency-stabilized UV laser having a tunable central offset frequency. The ULE cavity is used to stabilize the frequency of two seed lasers, an erbium-doped fiber laser at 1560.5 nm and an ytterbium-doped fiber laser at 1076.9 nm. The erbium-doped laser is directly locked to the ULE cavity, while the locking reference signal for the ytterbium-doped laser comes from the sum frequency beam at 637.2 nm. With a wideband waveguide-type EOPM for this ESB scheme, we obtain a continuous tuning range of 4 GHz in the UV at 318.6 nm, which is much wider than that reported previously [16]. The doubling cavity also remains locked to the 637.2 nm light. The laser system enables single-photon Rydberg excitation of cesium atoms from the 6S1/2 state to nP (n = 70–100) states.

### 2. Principle and experimental arrangement

As shown in figure 1, the EOPM2 is driven by a phase-modulated signal. The drive signal produced by a microwave function generator (MW-FG) has a carrier frequency of \( \omega_1 \) with a modulation depth of \( \beta_1 \), and is phase modulated by a radio frequency function generator (RF-FG2) at \( \omega_2 \) with a modulation depth of \( \beta_2 \). Therefore, the electric field of the laser exiting the EOPM2:

\[
E_{\text{ESB}} \approx E_0 \exp \{ i \omega t + \beta_1 \sin (\omega_1 t + \beta_2 \sin (\omega_2 t)) \},
\]

where \( E_0 \) is the amplitude of the incident laser at \( t = 0 \), and \( \omega \) is the angular frequency of the incoming beam. Using Bessel functions and ignoring contributions from high-order terms, we can expand this expression to first order in \( \beta \) (i = 1, 2).

\[
E_{\text{ESB}} \approx E_0 [ J_0(\beta) + 2i J_1(\beta) \sin (\Omega t + \beta_2 \sin (\Omega_2 t))] e^{i\omega t}
\approx E_0 [ J_0(\beta) e^{i\omega t} + E_0 J_1(\beta) \sin \Omega t]
\times [ J_0(\beta_2) e^{i(\omega_2 t + \Omega_2 t)} + J_1(\beta_2) e^{i(\omega_2 t + \Omega_2 t)}]
\times [ J_0(\beta_1) e^{i(\omega_1 t - \Omega_1 t)} - J_1(\beta_1) e^{i(\omega_1 t - \Omega_1 t)} - E_0 J_2(\beta_1) e^{i(\omega_1 t - \Omega_1 t)}]
\times [ J_0(\beta_2) e^{i(\omega_2 t - \Omega_2 t)} + J_1(\beta_2) e^{i(\omega_2 t - \Omega_2 t)} + E_0 J_2(\beta_2) e^{i(\omega_2 t - \Omega_2 t)}].
\]

Here, \( J_n(\beta) \) is the \( n \)-order Bessel function. The result of the cascaded phase modulation is to split the light into seven different frequency components: a carrier at \( \omega \), two sidebands at \( \omega \pm \Omega_1 \), and four sub-sidebands at \( \omega \pm \Omega_1 \pm \Omega_2 \). The two components at \( \omega + \Omega_1 \pm \Omega_2 \) (or \( \omega - \Omega_1 \pm \Omega_2 \)) have opposite phases. Therefore, when the sideband at \( \omega + \Omega_1 \) or \( \omega - \Omega_1 \) is locked to a fixed reference cavity, and the carrier and other sidebands are reflected, the ESB error signal is obtained by demodulating the detected signal with the modulation frequency at \( \Omega_2 \). The frequency discriminant of the locking technique is given by

\[
D_{\text{ESB}} = \frac{16FLE_0^2}{c} J_0(\beta_2) J_1(\beta_1) J_2(\beta_1),
\]

where \( c \) is the speed of light in vacuum, \( L \) is the optical cavity length, and \( F \) is the finesse defined by \( F = \text{FSR}/\Delta \nu \). Here, \( \Delta \nu \) is the full width at half maximum of the cavity signal. When one of the sidebands is near the resonance frequency of the cavity, such as \( \omega + \Omega_1 = 2\pi n \cdot \text{FSR} \), the carrier frequency of the laser can be tuned by adjusting \( \Omega_1 \).

A schematic of the experimental arrangement is shown in figure 2. Both infrared lasers (NKT Photonics) seed commercial wide gain bandwidth amplifiers, a 15 W erbium-doped fiber amplifier (EDFA) and a 10 W ytterbium-doped fiber amplifier (YDFA). The boosted laser beams are combined and passed through a 40 mm × 10 mm × 0.5 mm MgO-doped periodically poled lithium niobate crystal with a poling period of 11.80 \( \mu \)m. Previously [17], the crystal was temperature-stabilized at 154.0 \( ^\circ \)C to produce 8.75 W of 637.2 nm light using single-pass sum frequency generation. The 637.2 nm light is separated by a half-wave plate and a polarization beam splitter cube. The main portion of the 637.2 nm light is frequency doubled to 318.6 nm in a self-designed cavity using a 10 mm × 3 mm × 3 mm Brewster-cut \( \beta \)-BaB\(_2\)O\(_4\) (BBO) crystal. A small fraction of its output (\( \sim 3 \) mW) is phase modulated by a fiber-coupled waveguide-type EOPM2 (Jenoptik PM635), and then ESB locked to the ULE cavity to stabilize the 1076.9 nm laser.

A waveguide-type EOPM1 (EO-Space PM-055-10-PFA-PFA-UL) is placed between the distributed feedback ytterbium-doped fiber laser (DFB-EDFL) and EDFA to produce a set of sidebands at a modulation frequency of 12.6 MHz. The phase-modulated 1560.5 nm laser output is split and one beam is injected into the ULE cavity, and the response is monitored by PD1 in the reflected cavity signal. The driven
The signal of a function generator is phase shifted, and then mixed with the detected signal from PD1 to stabilize the 1560.5 nm laser frequency using the PDH technique [14].

The ULE reference cavity is a two-mirror spherical Fabry-Pérot cavity (AT Films) consisting of one plane-plane mirror and one plane-concave mirror with a 500 mm radius of curvature. Both mirrors and the cavity body consist of ULE glass. The mirrors are coated for high reflectivity at 1560.5 and 637.2 nm. The optical cavity is 47.6 mm long, resulting in a 3.145 GHz free spectral range (FSR). With a design similar to that described in [18], the optical cavity is housed in a thermal radiation shield inside a temperature-stabilized and ultra-high vacuum chamber to ensure a uniform temperature that is actively stabilized at the zero crossing of the cavity’s coefficient of thermal expansion of 11.13(56) °C. The pressure in the vacuum chamber is kept at ~3 × 10^{-9} Torr using an 8 L s^{-1} ion pump. According to the modulation sideband method, the finesse of the ULE cavity is 3.4(2) × 10^4 (FWHM ~ 92 kHz) at 1560.5 nm and 3.0(2) × 10^5 (FWHM ~ 105 kHz) at 637.2 nm.

### 3. Results and discussion

A tunable 318.6 nm laser is required to address various Rydberg states. By slowly changing the DFB-EDFL and DFB-YDFL temperatures over the range 20°C–50 °C, the 1560.5 and 1076.9 nm infrared seed lasers can be coarsely tuned over 145 and 202 GHz, respectively. This allows coarse tuning ranges of 347 and 694 GHz for the 637.2 and 318.6 nm lasers, respectively. When a tunable sideband of the 637.2 nm laser is locked to the ULE cavity, the offset between the target and reference frequencies depends on the frequency \( \Omega_{\text{gen}} \) generated by MW-FG (Agilent E8257C). The electronic signal is phase modulated at \( \Omega_2 \) to produce a pair of sidebands for ESB locking. The PD2 output signal is demodulated with \( \Omega_2 \) generated by RF-FG2 (Agilent 33250A) via a phase shift. To optimize the error signal, \( \Omega_2/2\pi \) is set to 2 MHz with a 10 dBm modulation amplitude. The ESB error signal from the mixer goes into a proportional-integral-differential controller after a 1.9 MHz low-pass filter, and then is fed back to the piezo-electric transducer (PZT) of the 1076.9 nm YDFL to produce frequency-stabilized tunable 637.2 nm laser. The
RF power consumptions of 14 and 10 dBm, respectively. The applied modulation frequency $\Omega_1/2\pi$ and $\Omega_2/2\pi$ are equal to 15 and 2 MHz with RF power consumptions of 14 and 10 dBm, respectively.

The blue curve represents the corresponding ESB error signal. Here, $\Omega_1/2\pi$ and $\Omega_2/2\pi$ are equal to 15 and 2 MHz with RF power consumptions of 14 and 10 dBm, respectively.

When the carrier frequency of the 1560.5 nm laser and the upper sideband of the 637.2 nm light are locked to the ULE cavity, we tune the carrier frequency of the 637.2 nm light by changing the phase modulation frequency $\Omega_1$. The frequency $\Omega_1/2\pi$ is automatically swept across $\sim2.4$ GHz over the range 1.1–3.5 GHz in a period of 1000 s. Thus, the frequency-tuning rate of the red light is 2.4 MHz s$^{-1}$. These parameter settings are established by the internal automatic sweep function of MW-FG. As shown in figure 6, if the tuning rate of the 637.2 nm laser frequency is up to hundreds of MHz s$^{-1}$, where the phase noise of the locking loops is very close to that of the loose lock, it will cause instability in the feedback loops. Therefore, to keep locking loops stable with low phase noise when the laser frequency is swept, a relatively low rate of frequency tuning (less than several MHz s$^{-1}$) should be used to address the desired state. The tuning range is characterized by a monitor cavity with a 487 MHz FSR. The 637.2 nm laser can be swept across more than four FSRs of the monitor cavity while maintaining lock. Meanwhile, the 318.6 nm UV laser is swept over eight FSRs of the monitor cavity with a FSR of $\sim500$ MHz, indicating that the continuous tuning ranges of the stable 637.2 and 318.6 nm lasers are over 1.95 and 4 GHz, respectively (see figure 5). The tuning range depends mainly on the gain-flattening region of the EOPM and the PZT-tunable range of the 1076.9 nm YDFL. In addition, it is also limited by the RF bandwidth of the EOPM.

When the 1560.5 nm laser is locked to the ULE cavity and the 1076.9 nm laser is free-running and swept by a triangular wave to address the desired Rydberg state, the 637.2 nm laser monitored by the ULE cavity varies about $\sim3.5$ GHz in a period of 1000 s. Thus, the frequency stability is estimated from the slope of the zero-crossing point of the ESB error signal. A 2000 s time trace of the PSDs above 17 kHz arises from the response of the error signal to high-frequency deviations beyond its linear region.

The modulation index of the EOPM2 depends mainly on the applied modulation frequency $\Omega_1$, even if the modulation amplitude is constant (figure 4). The modulation index decreases sharply at modulation frequencies less than 1.1 GHz. Therefore, when sweeping the frequency of the modulator, the gain of the error signal feedback to the PZT of the 1076.9 nm fiber laser is reduced to prevent the 637.2 nm laser frequency from being locked. To obtain a frequency-stabilized tunable 318.6 nm laser, the tuning range of the carrier frequency of the 637.2 nm laser is chosen in the range 1.1–3.5 GHz, which is in the gain-flattening region of the EOPM2 modulation index. In the locked condition, the upper bound of the 637.2 nm laser tuning range is limited by the 3.5 GHz piezo-tuning range of the 1076.9 nm YDFL.

When the carrier frequency of the 1560.5 nm laser and the upper sideband of the 637.2 nm light are locked to the ULE cavity, we tune the carrier frequency of the 637.2 nm light by changing the phase modulation frequency $\Omega_1$. The frequency $\Omega_1/2\pi$ is automatically swept across $\sim2.4$ GHz over the range 1.1–3.5 GHz in a period of 1000 s. Thus, the frequency-tuning rate of the red light is 2.4 MHz s$^{-1}$. These parameter settings are established by the internal automatic sweep function of MW-FG. As shown in figure 6, if the tuning rate of the 637.2 nm laser frequency is up to hundreds of MHz s$^{-1}$, where the phase noise of the locking loops is very close to that of the loose lock, it will cause instability in the feedback loops. Therefore, to keep locking loops stable with low phase noise when the laser frequency is swept, a relatively low rate of frequency tuning (less than several MHz s$^{-1}$) should be used to address the desired state. The tuning range is characterized by a monitor cavity with a 487 MHz FSR. The 637.2 nm laser can be swept across more than four FSRs of the monitor cavity while maintaining lock. Meanwhile, the 318.6 nm UV laser is swept over eight FSRs of the monitor cavity with a FSR of $\sim500$ MHz, indicating that the continuous tuning ranges of the stable 637.2 and 318.6 nm lasers are over 1.95 and 4 GHz, respectively (see figure 5). The tuning range depends mainly on the gain-flattening region of the EOPM and the PZT-tunable range of the 1076.9 nm YDFL. In addition, it is also limited by the RF bandwidth of the EOPM.

We characterize the laser frequency stabilization by comparing the power spectral densities (PSDs) of the closed-loop ESB error signal for two cases: a tight lock with optimized parameters and a loose lock where the stabilization is just sufficient to keep the laser frequency on the central slope of the error signal [19]. A nearly linear voltage response for frequency fluctuations is obtained. In figure 6, the experimental data are analyzed with a fast Fourier transform to extract the lower part of the frequency noise power spectrum (10Hz to 20 kHz). It can be seen that the phase noise has been reduced by more than 30 dB within a 17 kHz bandwidth from the crossing point of the two curves. The difference between the PSDs above 17 kHz arises from the response of the error signal to high-frequency deviations beyond its linear region.

If the frequency difference between the sidebands and sub-sidebands of the transmission spectrum is used as a ruler, we lock the upper sideband of the carrier frequency to the zero-crossing point of the two curves. The difference between the PSDs above 17 kHz arises from the response of the error signal to high-frequency deviations beyond its linear region.

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The modulation index dependence of EOPM2 applied modulation frequency. Here the EOPM2 (Jenoptik PM635) is phase modulated by MW-FG (Agilent E8257C).

ESB error signal and the transmitted signal of the phase-modulated 637.2 nm light incident on the ULE cavity are shown in figure 3.

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shown in the inset of figure 7. The relative frequency fluctuations of the locked 637.2 nm laser are below 8 kHz. Figure 7 shows the relative Allan deviation of the 637.2 nm light. Its relative frequency instability is less than $1.5 \times 10^{-11}$ for interrogation times between 0.1 and over 500 s, corresponding to a frequency deviation of $\sim 7$ kHz. Thus, the relative frequency deviation of the 318.6 nm laser is estimated to be less than 15 kHz. It should be emphasized that the relative Allan deviation is derived from the error signal. The measurement is relatively insensitive to frequency variations introduced by tiny variations in cavity length, so it only represents a lower limit of the frequency instability. The Allan deviation should be measured in more detail by the frequency beating of two identical stable laser systems, or optical frequency comb technology, which more accurately reflects laser frequency stability.

Due to the residual amplitude modulation in the EOPM, we can either control the modulator temperature or introduce an intense optical beam to suppress it [20]. To further improve the stability, a fast feedback loop needs to be constructed by inserting an AOM in the 637.2 nm light path [21]. The resulting frequency perturbations of the red light are corrected by feedback on the frequency and amplitude of the AOM, a PZT inside the 1076.9 nm YDFL, and its temperature. In addition, novel materials play an important role in reducing the phase noise of the generated pulse [22, 23], which is possible to improve our fiber laser performance. However, the frequency stability is still sufficient for single-photon Rydberg excitations of cesium atoms [8].

4. Conclusion

In summary, we have demonstrated a continuously tunable frequency-stabilized 318.6 nm UV laser system. A high-finesse ULE optical cavity inside a temperature-stabilized
ultra-high vacuum chamber is used as frequency reference to stabilize 1560.5 and 1076.9 nm seed fiber lasers. Based on a commercial wideband waveguide-type EOPM and the ESB locking technique, the 637.2 nm laser can be continuously tuned over a range of 1.95 GHz while in lock. Meanwhile, the high-stability 318.6 nm UV laser can be continuously tuned over 4 GHz. Further improvements in the tuning range can be achieved by increasing the RF bandwidth of the EOPM, using a widely tunable PZT, and designing the automatic gain control circuit to adjust servo loop parameters. The frequency-stabilized UV laser system with a tunable central frequency is important for cesium experiments, including single-photon Rydberg excitation and high precision spectroscopy of Rydberg excitation and high precision spectroscopy of Rydberg atoms and their interactions.

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